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Axions in physics and astrophysics

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June 13, 2018

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Overview

- 1 Not Only Theoretical motivation
 - QCD Lagrangian
 - Neutron electric dipole momentum
- 2 ALPs in cosmology
 - ALPs as a dark matter
 - Cosmological bounds
- 3 Indirect ALPs detection
 - Ways to detect
- 4 Astrophysical constraints on ALPs
 - Most interesting astrophysical constrains on ALPs
 - SN 1987A
 - TeV transparency of the Universe
 - Photon-Axion conversion in the clusters of galaxies

5 What's next?

Most generic QCD Lagrangian

A "text-book" QCD Lagrangian is CP-invariant

$$\begin{split} \mathcal{L}_{QCD} &= -\frac{1}{4} G_{\mu\nu} G^{\mu\nu} + i \bar{\psi} D_{\mu} \gamma^{\mu} \psi + \bar{\psi} M \psi \\ G^{a}_{\mu\nu} &= \partial_{\mu} A^{a}_{\nu} - \partial_{\nu} A^{a}_{\mu} + g f^{abc} A^{b}_{\mu} A^{c}_{\nu} - \text{gluon field strength tensor} \\ A - \text{gluons} \\ \psi - \text{quarks} \end{split}$$

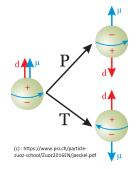
But it is not the most general form of 4-scalar built from G. θ -term can be added:

$$\mathcal{L}_{QCD}^{VCP} \sim \mathcal{L}_{QCD} + \theta \varepsilon^{\mu\nu\alpha\beta} \mathcal{G}_{\mu\nu} \mathcal{G}_{\alpha\beta}$$

 θ - term behaves similar to $\vec{E} \cdot \vec{B}$ in QED. Due to θ -term \mathcal{L}_{QCD}^{VCP} is **CP-violating** and **T-violating**

CP-violation... So what?

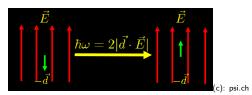
- Any quark-built particle with electric dipole and magnetic momenta violates T (and CP) symmetries (ok, we expect this with L^{VCP}_{QCD})
- Neutron consists of (*udd*) quarks and has very well measured magnetic momentum



- Neutron electric dipole momentum (nEDM): $d \sim e \times \text{length} \times \theta \sim e \times \frac{m_q}{\Lambda_{QCD}^2} \times \theta \sim (3..30) \cdot 10^{-16} \times \theta[e \cdot \text{cm}]$
- Very similar result from naive estimation neutron is a particle consisting of 3 charged quarks.

Strong CP problem





 $n\mathsf{EDM}=0-\mathsf{CP}\ holds\quad;\quad n\mathsf{EDM}>0-\mathsf{CP}\ violates$

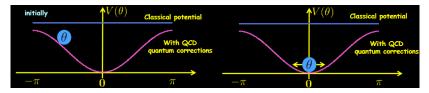
Strong CP problem:

- Current experimental limit $d < 3 \cdot 10^{-26} \ [e \cdot \text{ cm}]$
- CP holds. Tiny, consistent with zero θ . Fine tuning!

Zero θ is just a coincidence? Don't think so...

- CP is *not* fundamental symmetry and is known to be violated in weak interactions (kaon, b-meson decays)
- Cosmological requirements for CP-violation (baryogenesis)

Introducing axion/ALP



(c):https://www.psi.ch/particle-zuoz-school/Zuoz2016EN/jaeckel.pdf

Solution for strong CP problem – make θ dynamical variable.

- Produced with non-zero value after some time it can relax to zero. New particle gained – <u>axion</u>! (~ 1977 Peccei,Quinn+)
- Historically axion particle with specific $m_a g_a$ relation
- More general type of particles ALPs (axion-like particles)
- ALPs are (generally) light, since CP violation is tiny
- ALPs are WISPs Weakly Interacting Sub-eV Particles

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ALPs as a dark matter

Can ALPs play a role in cosmology/astrophysics ?

Successful dark matter candidate:

• Massive particles \rightarrow ok!

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- Massive particles \rightarrow ok!
- Non-relativistic (CDM/WDM) \rightarrow ok! (not that obvious!)
- Produced in Early Universe \rightarrow ok! (not that obvious!)

 Despite of low mass ALPs indeed can be produced in early universe out of thermal equilibrium via vacuum realignment or topological production mechanisms (see e.g. 1712.03018 and 0904.3346)

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- (Optional) Other than DM motivation/manifestation \rightarrow ok!

ALPs as a dark matter

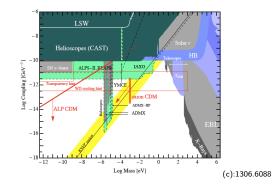
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- Indeed, ALPs are good dark matter candidate!
- Note: ALPs can exist and *not* to be a dark matter!

Cosmological bounds on ALPs parameters



- Arise from comparison of (model-dependent) amount of ALPs produced in Early-Universe to Ω_{CDM}
- Not-unavoidable, but rather order-of-magnitude estimation!
- Functional form

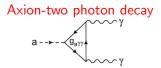
$$\Omega_a \sim m_a^{1/2} g_a^{-2} = const$$
 ; see 1210.5081

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Ways to detect

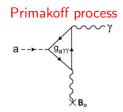
Axions are coupled to "direct observables" – photons via triangle diagrams.



One-to-one correspondence between photon energy E_{γ} and axion mass m_a :

$$E_{\gamma} = m_a/2$$

$$\Gamma = 7.6 \cdot 10^{-26} \left(\frac{g_a}{10^{-10} GeV^{-1}} \right)^2 \left(\frac{m_a}{1eV} \right)^3 s^{-1}$$



In a presence of (electro) magnetic field axion can be converted to a photon and v.v: photon-axion oscillations

Primakoff process

• General case:

$$(E - i\partial_z - M)\vec{A} = 0; \vec{A} = \begin{pmatrix} A_x \\ A_y \\ a \end{pmatrix}$$

M - real symmetric 3x3 matrix; $M = M(m_a, g_a, \vec{B}(z), n_e(z))$ $P_{\gamma \to a} \propto |AA^{\dagger}|^2 = P_{\gamma \to a}(E_{\gamma}, s, \vec{B}, n_e, m_a, g_a)$

• Formal solution for \vec{A} after traveling distance *s* with magnetic field domain size δz :

$$\vec{A} = \prod_{j=1}^{s/\delta z} \exp\left(iM(z_j)\delta z\right) \vec{A}(0)$$

• Exact $P_{\gamma
ightarrow a}$ known for 1 and ∞ magnetic domains.

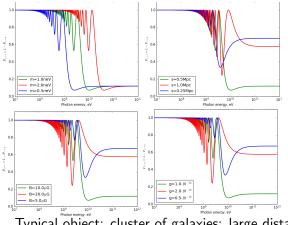
Primakoff process

- For fixed m_a P_{γ→a} has local maxima at certain photon energies E_{γ,max}(B, n_e) – resonant axion production. At these energies we expect features in spectra of astrophysical objects.
- NO one-to-one correspondence: m_a can be very different from $E_{\gamma,\max}$
- For objects with known \vec{B} and n_e we can probe axion mass range for which $E_{\gamma,\max}$ are located in "preferable" band (keV, GeV, TeV).
- (Non)Detection of such spectral features can allow constrain $m_a g_a$ parameter space or find axions.

Examples of $P_{\gamma \rightarrow a}$ in astrophysics

For most of astrophysical objects one of two regimes works:

• Magnetic field dominating regime:



Oscillations' energy strongly correlate with m_a and strength – with g_a and B. Generally strong constraints on $m_a - g_a$, but in narrow axion mass range

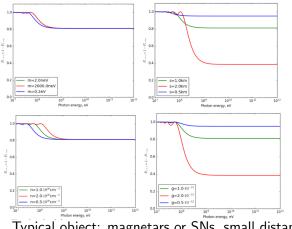
Typical object: cluster of galaxies: large distance, low B, $n_e \sim 0$

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Examples of $P_{\gamma \rightarrow a}$ in astrophysics

For most of astrophysical objects one of two regimes works:

• Plasma dominating regime:



Oscillations' energy strongly correlate with n_e and strength – with g_a and B. $P_{\gamma \rightarrow a}$ almost does not depend on m_a ! Weak constraints but in broad $m_a - g_a$ range.

Typical object: magnetars or SNs, small distance, high B and n_e

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Astrophysical ways to constrain ALPs

Spectral line from axion-two photon decay can be used for searches of ALP-dark matter.

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- Spectral line from axion-two photon decay can be used for searches of ALP-dark matter.
 - Remember: ALPs are not necessary dark matter!
 - Can be complicate for low (radio–optics) energies and axion mass (remember $E_{\gamma}=m_a/2$)

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 - Remember: m_a can be different from spectral feature energy!
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 - (non)transparency of the medium to photons of some (${\sim} {\rm TeV})$ energies

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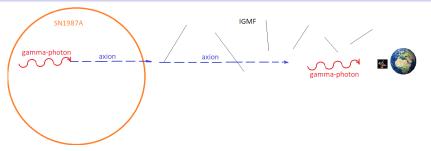
• combination of both approaches

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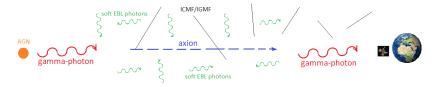
Changes in evolution of some objects (white dwarf cooling/massive stars evolution)

Most famous constraint: SN 1987A



- Galactic supernova: first detection of astrophysical neutrino
- Neutrino signal implies high intrinsic photon flux expected high ALP flux due to conversion in SN (dense plasma)
- ALPs travel through inter-galactic magnetic field to the Earth
- ALPs can be converted to photons during the travel
- no γ -rays have been seen (by SMM/GRS built for the Sun observations)

Most interesting constraint: lower-limit from transparency



- VHE (> 1 TeV) photons are effectively absorbed for distant sources (pair production with soft EBL photons)
- Some TeV signal/spectral hardening seen in distant AGNs
- Understood if some distance photons travel as ALPs
- Lower limit: g_a > ...
- A lot of assumptions: EBL spectral/density model ; ICMF/IGMF ; no intrinsic hardening of AGN spectra, etc...

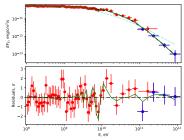




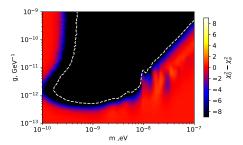
- Bright gamma/X-ray source (AGN) located within cluster of galaxies
- Spectral features ("absorption lines") can be produced by resonant photon-ALP conversion at certain energies
- Intrinsic spectrum believed to be featureless (powerlaw)
- $\bullet\,$ Low-background objects at $\gtrsim 10~{\rm keV}$ energies
- Key ingredient magnetic field profile in a cluster



Photon-axion in Perseus cluster/NGC 1275



 ${\sf Fermi/LAT}{+}{\sf MAGIC}$ spectrum of NGC 1275 with the best-fit axion-absorption model



95% c.l. limits on $m_a - g_a$ parameter space derived from NGC 1275 spectrum; (c):1805.04388

- Perseus cluster is one of the closest to us cluster of galaxies. NGC 1275 (central AGN) is a bright GeV-TeV source
- Fermi/LAT+MAGIC (0.01–1 TeV) data available
- Smooth spectrum with marginal ($\sim 2.5\sigma$ oscillations)
- Non-detection of oscillations \rightarrow constraints on $m_a g_a$

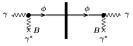
Few words on direct-detection experiments

Solarscope



- X-ray photons converted to axions inside of the Sun
- Converted back in a strong magnetic field
- Experiments: CAST, IAXO

Light through the wall



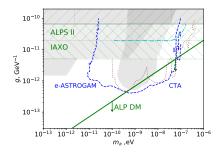
- (c): https://alps.desy.de/
 - Forth-and-back conversion in a magnetic field
 - $\gamma \rightarrow a$
 - wall
 - $a \rightarrow \gamma$
 - Experiments: ALPS, ALPS II

Resonant cavity



- DM axions resonantly convert to γ in a cavity with magnetic field
- Experiments: ADMX. ADMX-HF ୬**୯**୯ 19/22

What's next? Instead of conclusions



Next instruments within 10 years – e-ASTROGAM and CTA. Much broader energy range+higher data quality=better and broader exclusion region for ALP parameters (or detection?)

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- Sensitivity of indirect searches is comparable to the sensitivity of the near future direct-detection experiments (IAXO,ALPS II), although in narrower mass range.
- See also arXiv:1805.04388

Really last slide!

 Any object with high magnetic field/plasma density can be used to probe "ALP-effects"

- Key element knowledge of parameters (\vec{B}, n_e) of the analyzed object
- Stacking analysis can significantly increase sensitivity
- This is not the end of the story with NGC 1275...



The End!



"One needs a particle to clean up a problem..." -Frank Wilczek