

## Advanced Econometrics - 1<sup>st</sup> assignment sheet

### Task 1

Show that these functions are Kernels:

$$K(\psi) = (2\pi)^{-1/2} \exp(-0.5\psi^2) \quad (1)$$

$$K(\psi) = \mathbf{1}\{|\psi| \leq 1\} 3/4 \cdot \{1 - \psi^2\} \quad (2)$$

$$K(\psi) = \mathbf{1}\{|\psi| \leq 1/2\} \quad (3)$$

### Solution

Kernel 1: We know that the normal density integrates up to one.

Kernel 2:

$$\begin{aligned} \int_{-\infty}^{+\infty} K(\psi) &= \int_{-1}^{+1} K(\psi) \\ &= \int_{-1}^{+1} 3/4 \cdot (1 - \psi^2) d\psi \\ &= \frac{3}{4} \left[ \psi - \frac{\psi^3}{3} \right]_{-1}^1 \\ &= \frac{3}{4} \left( \frac{2}{3} + \frac{2}{3} \right) \\ &= 1 \end{aligned}$$

Kernel 3:

$$\begin{aligned} \int_{-\infty}^{+\infty} K(\psi) &= \int_{-1/2}^{+1/2} K(\psi) \\ &= \int_{-1/2}^{+1/2} 1 d\psi \\ &= [\psi]_{-1/2}^{1/2} \\ &= \left( \frac{1}{2} + \frac{1}{2} \right) \\ &= 1 \end{aligned}$$

### Task 2

Show that the relation

$$\int_{-\infty}^{+\infty} \psi^k K(\psi) d\psi = 0 \quad (4)$$

holds, for  $k = 1, 3, 5, \dots$

*Solution*

Fist we take a look at the case where  $k = 1$

$$\begin{aligned} \int_{-\infty}^{+\infty} \psi K(\psi) d\psi &= \int_{-\infty}^0 \psi K(\psi) d\psi + \int_0^{+\infty} \psi K(\psi) d\psi \\ &= - \int_0^{-\infty} \psi K(\psi) d\psi + \int_0^{+\infty} \psi K(\psi) d\psi \\ &= - \int_0^{-\infty} \int_0^\psi 1 dt K(\psi) d\psi + \int_0^{+\infty} \int_0^\psi 1 dt K(\psi) d\psi \\ &= - \int_0^{-\infty} \int_t^{-\infty} K(\psi) d\psi dt + \int_0^{+\infty} \int_t^{\infty} K(\psi) d\psi dt \\ &= - \int_0^{\infty} \int_t^{\infty} K(\psi) d\psi dt + \int_0^{+\infty} \int_t^{\infty} K(\psi) d\psi dt = 0 \end{aligned}$$

The last step is done, using the symmetry property.

*Task 3*

Proof the following Collary:

**Corollary 1** *Let assumptions DE I-IV hold. Then*

$$plim \hat{f}_x(x_0) \rightarrow f_x(x_0) \quad (5)$$

*Solution*

Proof of Collary 1:

At first: It can be show that if a random variable  $x_n$  converges in mean square to a constant  $c$ , then it converges in probability to  $c$  (See Greene 4th ed<sup>1</sup> p. 110, for example). This can be shown by using *Chebychev's Inequality*. Further it can be shown that if

$$\lim_{n \rightarrow \infty} E[x_n] = c \text{ and } \lim_{n \rightarrow \infty} \text{Var}[x_n] = 0,$$

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<sup>1</sup>In the 5th ed p.897/898

that  $x_n$  converges in mean square and hence in probability to  $c$  (See also Greene p. 110/111).

So it is sufficient for showing  $plim \hat{f}_x(x_0) \rightarrow f_x(x_0)$  that  $\mathbf{E}[\hat{f}_x(x_0)] = f_x(x_0)$  and  $\mathbf{Var}[\hat{f}_x(x_0)] = 0$ .

$$\begin{aligned}
\mathbf{E}[\hat{f}_x(x_0)] &= \mathbf{E}\left[\frac{1}{nh}\sum_{i=1}^n K\left(\frac{x_i - x_0}{h}\right)\right] \\
&\stackrel{iid}{=} \frac{1}{h}\mathbf{E}\left[K\left(\frac{x_i - x_0}{h}\right)\right] \\
&= \frac{1}{h}\int_{-\infty}^{+\infty} K\left(\frac{\xi - x_0}{h}\right) f_x(\xi) d\xi \\
&= \int_{-\infty}^{+\infty} K(\psi) f_x(h\psi + x_0) d\psi \\
&= \int_{-\infty}^{+\infty} K(\psi) \left(f_x(x_0) + h\psi\partial_x f_x(x_0) + \frac{h^2\psi^2}{2}\partial_x^2 f_x(x_0) + \frac{h^3\psi^3}{6}\partial_x^3 f_x(x_0)\right) d\psi \\
&= f_x(x_0) + h\partial_x^2 f_x(x_0) \int_{-\infty}^{+\infty} K(\psi)\psi^2 d\psi + \frac{h^3}{6} \int_{-\infty}^{+\infty} \partial_x^3 f_x(x_r)\psi^3 K(\psi) d\psi \\
&= f_x(x_0) + \frac{h^2}{2}\partial_x^2 f_x(x_0)\mu_2 + \frac{h^3}{6} \int_{-\infty}^{+\infty} \partial_x^3 f_x(x_r)\psi^3 K(\psi) d\psi \\
&= f_x(x_0) + \frac{h^2}{2}\partial_x^2 f_x(x_0)\mu_2 + o(h^2)
\end{aligned}$$

The last step can be retraced by recognizing that:

$$\begin{aligned}
\frac{h^3}{6} \int_{-\infty}^{+\infty} \partial_x^3 f_x(x_r)\psi^3 K(\psi) d\psi &\leq \left|\frac{h^3}{6} \int_{-\infty}^{+\infty} \partial_x^3 f_x(x_r)\psi^3 K(\psi) d\psi\right| \\
&\leq \frac{h^3}{6} \int_{-\infty}^{+\infty} |\partial_x^3 f_x(x_r)\psi^3 K(\psi)| d\psi \\
&\leq \frac{h^3}{6} \int_{-\infty}^{+\infty} |\partial_x^3 f_x(x_r)||\psi^3|K(\psi) d\psi \\
&\leq \frac{h^3}{6} b \int_{-\infty}^{+\infty} |\psi^3|K(\psi) d\psi \\
&\leq \frac{h^3}{6} bc
\end{aligned}$$

since  $|\partial_x^3 f_x(x_r)|$  is assumed to be bounded and  $\int_{-\infty}^{+\infty} |\psi^3|K(\psi) d\psi$  is also bounded.

We show now that the Variance vanishes for  $n \rightarrow \infty$ .

At fist the Variance can be written as

$$\begin{aligned}
\text{Var}[\hat{f}_x(x_0)] &= \text{Var}\left[\frac{1}{nh}\sum_{i=1}^n K\left(\frac{x_i - x_0}{h}\right)\right] \\
&= \frac{1}{(nh)^2}n\text{Var}\left[K\left(\frac{x_i - x_0}{h}\right)\right] \\
&= \frac{1}{nh^2}\left(\mathbb{E}\left[K\left(\frac{x_i - x_0}{h}\right)^2\right] - \mathbb{E}\left[K\left(\frac{x_i - x_0}{h}\right)\right]^2\right).
\end{aligned}$$

Therefore we have to calculate  $\mathbb{E}\left[K\left(\frac{x_i - x_0}{h}\right)^2\right]$ .

$$\begin{aligned}
\mathbb{E}\left[K\left(\frac{x_i - x_0}{h}\right)^2\right] &= \int_{-\infty}^{\infty} K\left(\frac{\xi - x_0}{h}\right)^2 f_x(\xi) d\xi \\
&= \int_{-\infty}^{\infty} K(\psi)^2 f_x(\psi h + x_0) h d\psi \\
&= h \int_{-\infty}^{\infty} K(\psi)^2 (f_x(x_0) + h^2 \psi \partial_x f_x(x_r)) d\psi \\
&= h f_x(x_0) \int_{-\infty}^{\infty} K(\psi)^2 d\psi + h^2 \int_{-\infty}^{\infty} \psi K(\psi)^2 \partial_x f_x(x_r) d\psi \\
&= h f_x(x_0) c + o(h)
\end{aligned}$$

The last equality holds because of:

$$\begin{aligned}
h^2 \int_{-\infty}^{\infty} \psi K(\psi)^2 \partial_x f_x(x_r) d\psi &\leq |h^2 \int_{-\infty}^{\infty} \psi K(\psi)^2 \partial_x f_x(x_r) d\psi| \\
&\leq h^2 \int_{-\infty}^{\infty} |\psi K(\psi)^2| |\partial_x f_x(x_r)| d\psi \\
&\leq h^2 b \int_{-\infty}^{\infty} K(\psi)^2 |\psi| d\psi \\
&\leq h^2 b c_2 = O(h^2) = o(h)
\end{aligned}$$

Now, the variance can be written as

$$\begin{aligned}
\text{Var}[\hat{f}_x(x_0)] &= \frac{1}{nh^2}(h f_x(x_0) c + O(h^2)) + \frac{1}{nh^2}\left(h f_x(x_0) + \frac{h^3}{2} \partial_x^2 f_x(x_0) \mu_2 + O(h^4)\right)^2 \\
&= \frac{1}{nh} f_x(x_0) c + O(n^{-1}) + \frac{f_x(x_0)}{n} + \frac{h^2}{n} f_x(x_0) \partial_x^2 f_x(x_0) \mu_2 \\
&\quad + \frac{h^4}{4n} (\partial_x^2 f_x(x_0))^2 \mu_2^2 + 2 \frac{1}{nh} f_x(x_0) \cdot O(h^4) + \frac{h}{n} \partial_x^2 f_x(x_0) \mu_2 O(h^4) + O(h^6/n) \\
&= O((nh)^{-1}) + O(n^{-1}) + O(n^{-1}) + O(h^2/n) + O(h^4/n) + O(h^3/n) + O(h^5/n) + O(h^6/n) \\
&= O((nh)^{-1}) = o(1)
\end{aligned}$$

*Task 4*

Show that the following theorem holds:

**Theorem 1** *Under ADE I', ADE II, ADE III' and ADE IV we obtain*

$$E[\hat{f}_x(x_0)] = f_x(x_0) + \frac{h^r}{r} \mu_r \partial_x^r f_x(x_0) + o(h^r) \quad (6)$$

$$\text{Var}[\hat{f}_x(x_0)] = (nh)^{-1} \kappa_0 f_x(x_0) + o((nh)^{-1}) \quad (7)$$

*Solution*

$$\begin{aligned} E[\hat{f}_x(x_0)] &= E \left[ \frac{1}{hn} \sum_{i=1}^n K \left( \frac{x_i - x_0}{h} \right) \right] \\ &= \frac{1}{hn} n \int_{-\infty}^{\infty} K \left( \frac{\xi - x_0}{h} \right) f_x(x_0) d\xi \\ &= \frac{1}{h} \int_{-\infty}^{\infty} K(\psi) f_x(h\psi + x_0) h d\psi \\ &= \int_{-\infty}^{\infty} K(\psi) \left( f_x(x_0) + \psi h \partial_x f_x(x_0) + \frac{\psi^2 h^2}{2} \partial_x^2 f_x(x_0) + \dots \right. \\ &\quad \left. + \frac{\psi^r h^r}{r!} \partial_x^r f_x(x_0) \frac{\psi^{r+1} h^{r+1}}{(r+1)!} \partial_x^{r+1} f_x(x_{rest}) \right) d\psi \\ &= f_x(x_0) \int_{-\infty}^{\infty} K(\psi) d\psi + h \partial_x f_x(x_0) \int_{-\infty}^{\infty} \psi K(\psi) d\psi + \dots \\ &\quad + \frac{h^r}{r!} \partial_x^r f_x(x_0) \int_{-\infty}^{\infty} K(\psi) \psi^r d\psi + \frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} K(\psi) \psi^{r+1} \partial_x^{r+1} f_x(x_{rest}) d\psi \\ &= f_x(x_0) + \frac{h^r}{r!} \partial_x^r f_x(x_0) \mu_r + \frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} K(\psi) \psi^{r+1} \partial_x^{r+1} f_x(x_{rest}) d\psi \\ &= f_x(x_0) + \frac{h^r}{r!} \partial_x^r f_x(x_0) \mu_r + o(h^r) \end{aligned}$$

To see why the last equation holds, we can use similar arguments like in the task before:

$$\begin{aligned}
\frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} K(\psi) \psi^{r+1} \partial_x^{r+1} f_x(x_{rest}) d\psi &\leq \left| \frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} K(\psi) \psi^{r+1} \partial_x^{r+1} f_x(x_{rest}) d\psi \right| \\
&\leq \frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} |K(\psi) \psi^{r+1} \partial_x^{r+1} f_x(x_{rest})| d\psi \\
&\leq \frac{h^{r+1}}{(r+1)!} \int_{-\infty}^{\infty} K(\psi) |\psi^{r+1}| |\partial_x^{r+1} f_x(x_{rest})| d\psi \\
&\leq \frac{h^{r+1}}{(r+1)!} b \int_{-\infty}^{\infty} K(\psi) |\psi^{r+1}| d\psi \\
&\leq \frac{h^{r+1}}{(r+1)!} bc = O(h^{r+1}) = o(h^r)
\end{aligned}$$

Since the  $r+1^{th}$  derivative of  $f_x(x)$  is assumed to be bounded and the same holds for the reminding integral term.

Variance is the same as in the case of the usual kernels.