NAPL DISSOLUTION1

1. INTRODUCTION

Already at the beginning of contaminated site investigation and remediation it became obvious how challenging it is to decontaminate a polluted site within a "short" time frame (< decades). At many contaminated sites, remediation could not be achieved within a decade (Travis and Doty, 1990; Travis and Macinnis, 1992). Meanwhile it is clear that remediation often takes many decades (if not centuries). Besides slow diffusion over large distances, slow dissolution kinetics of NAPLs are responsible for the persistence of subsurface contaminations. The dissolution rate of NAPL depends on their surface to volume ratio (blobs or ganglia or pools) since it is limited by mass transfer across the NAPL/water interface. Large coherent volumes of NAPL (oil layers, DNAPL pools) take long times to dissolve as summarized for different scenarios in Table 1.

Tab. 1: Time scales for contaminant removal with remediation options (from Eberhard and Grathwohl, 2002, Grathwohl, 1998: Groundwater Quality Conference)

* this as for excavations requires the knowledge of the exact location of the NAPL source

¹ Parts of this chapter come from the book: Diffusion in Natural Porous Media: Contaminant Transport, Sorption/Desorption and Dissolution Kinetics. Kluwer Academic Publishers, Boston, 224 p. (ISBN 0-7923-8102-5); Peter Grathwohl (1998)

2. FILM DIFFUSION

2.1 BOUNDARY LAYERS – SHERWOOD NUMBERS

The dissolution rate F_b of NAPL trapped as blobs or ganglia (or pools) in a porous medium (e.g., an aquifer) is limited by diffusion across the NAPL/water interface through an aqueous boundary layer of thickness δ (Fig. 2.1a; Fick's 1st law):

$$
F_b = \frac{D_{aq}}{\delta} A (C_o - C)
$$
 (2.1)

 D_{aq} is the aqueous diffusion coefficient. *C* and C_o denote the solute concentrations in the mobile (i.e., aqueous) phase and at the interface, respectively. For pure organic liquids and water, C_o is the water solubility (for mixtures see Raoult's law). The ratio D_{aq}/δ is the mass transfer coefficient k [e.g., in m s -1]. The overall dissolution rate depends on the surface area (A) of the NAPL available for mass transfer. Often specific surface areas are employed which for example yield the flux per unit volume porous media. For spherical blobs trapped per unit volume in the porous medium the specific surface area then is:

$$
A_o = \frac{3 \theta_o}{r_b} \tag{2.2}
$$

 A_0 is the area of the organic liquid (NAPL) per unit volume porous media (m² m⁻³ = m⁻¹). θ_0 denotes the NAPL occupied porosity per unit volume of porous medium and r_b the blob radius. θ_o can be calculated from the NAPL saturation (S°) and the porosity $(\theta_o = S^{\circ} n)$. For $S^{\circ} = 0.05$ (5%), a porosity of 30% and an effective "blob" radius of 1 mm, the specific interfacial area (A_0) is 45 m² m⁻³. Specific surface areas may be normalized to different reference volumes (pore volume or NAPL volume) and care should be taken to select the right one. In some cases (e.g., mineral dissolution), the overall surface area may not be completely accessible for mass transfer into the flowing water and reduced effective interfacial areas are sometimes used.

Fig. 2.1: Simplified scheme of dissolution of a NAPL blob (shaded - yellow) by film diffusion (*a*: grain radius; r_b : radius of blob trapped in pore; δ : effective thickness of an aqueous boundary layer "film")

Since the film thickness depends on the flow velocity and is not explicitly known, *k* is calculated from empirical relationships, which use a dimensionless constant, the Sherwood number. The Sherwood number *Sh* is known from various applications in chemical engineering, and for porous media is defined as:

$$
Sh = \frac{k \, d}{D_{aq}} = \frac{d}{\delta} \tag{2.3}
$$

d denotes here the characteristic length in the porous medium, e.g., the grain size. According to Eq. 2.3, *Sh* may be interpreted as the ratio between grain diameter and film thickness δ ($k = D_{aq}/\delta$). *Sh* depends on flow velocity, porosity, and viscosity of water. Various empirical relationships have been developed for the calculation of *Sh*, e.g., in terms of other dimensionless numbers such as the Schmidt number *Sc* and the Reynolds number *Re*. Fitzer et al. (1995) report the following correlation for laminar flow in packed beds:

$$
Sh = 1.9 Sc^{1/3}Re^{1/2}
$$
 (2.4)

The Schmidt number *Sc* denotes the ratio between the kinematic viscosity (μ) of water and the aqueous diffusion coefficient $(Sc = \mu / D_{aq})$. In groundwater, *Sc* is for many hydrophobic organic compounds is almost a constant (e.g., *Sc* = 2600 at 10°C; μ = 1.3 x 10⁻⁶ m² s⁻¹; D_{aq} = ca. 0.5 x 10⁻⁹ m² s⁻¹). The Reynolds number in porous media is based on the grain diameter (*d*), the flow velocity (v_a), and the viscosity μ ($R_e = d v_a / \mu$). In natural porous media (aquifers), *Re* is in most cases less than 1 (e.g.: $v_a = 2$ m d⁻¹ and $d = 2$ mm \Rightarrow $Re = 0.36$ and *Sh* according to Eq. 2.4 approx. 5; in this case δ would be about 1/5 of the grain diameter as shown in Fig. 2.1). Other correlations were investigated for the dissolution of solvents in sands (e.g., Miller et al., 1990; Powers et al., 1991). Powers et al. (1994), for example, found the following relationship in column experiments for the dissolution of naphthalene spheres:

$$
Sh = 36.8 \, Re^{0.654} \tag{2.5}
$$

In this case, the Schmidt number is not used explicitly (for $Re = 0.036$, this correlation results in *Sh* $= 4.2$, which compares reasonably well with the example above, Eq. 2.4). Note, that a large number of theoretical and empirical Sherwood relationships exist in literature. Liu et al. (2014) reports an elegant correlation for Sherwood numbers for the dissolution of NAPL residuals in porous media, which is only based on the Peclet number $(=$ the product of *Re* and *Sc*) as shown in Fig 2.2. This applies if the viscosity of the fluid does not play a major role, and if flow conditions are laminar (typically for groundwater flow). Also, notice that we get a similar solution for the dissolution of DNAPL pools, which is based on boundary layer theory (and some simplifying assumptions).

Fig. 2.2: Relationship between Sherwood number (*Sh*) and Peclet number (*Pe*) for mass transfer of compounds with different molecular diffusion coefficients in water (from Liu et al., 2014). *Sh* = *k* d/D_{aq} and *Pe* = v_a d/D_{aq} .

For the calculation of the dissolution rates, the specific surface area of the NAPL (A_o) is required, which, however, is not well known and difficult to determine. Therefore, a modified Sherwood Sh' number was introduced, which also accounts for A_0 :

$$
Sh' = \frac{kA_o d^2}{D_{aq}}\tag{2.6}
$$

Miller et al. (1990) found the following empirical correlation between *Re* and *Sh'* for the dissolution of NAPL in column experiments:

$$
Sh' = 12 \, Re^{0.75} \theta^{0.6} Sc^{0.5} \tag{2.7}
$$

Using a value of 2 600 for *Sc* yields $Sh' = 612$ $Re^{0.75}$ $\theta^{0.6}$. Imhoff et al. (1993) report the following correlation:

$$
Sh' = 150 \, Re^{0.87} \theta^{0.79} \tag{2.8}
$$

Further and similar empirical correlations are reported in Mayer and Miller (1996) and Imhoff and Miller (1996). It should be noted that these relationships were measured for specific porous media and NAPLs and are not necessarily valid for other systems. However, as shown in Fig. 2.4 mass transfer during the dissolution of NAPL droplets in many cases is fast, and local equilibrium conditions can be assumed at the field scale.

2.1 LENGTHS OF MASS TRANSFER ZONES

Whether mass transfer limitations must be considered or not depends on the degree of equilibrium achieved if water flows through a mass transfer zone (e.g., consisting of interfaces water / residual NAPL or just water/solids, e.g., during the dissolution of minerals). For external mass transfer resistance (e.g., film diffusion in the aqueous boundary layer surrounding NAPL droplets or crystallites) Fick's first law applies (extended from Eq. 2.1):

$$
F' = \frac{D_{aq}}{\delta} A_o \left(C_{w,eq} - C_w \right) \tag{2.9}
$$

which simply states that mass transfer occurs by molecular diffusion across a specific interfacial area A_0 and a characteristic distance δ (film thickness). *F'* denotes the flux per unit volume of porous medium. $C_{w,eq}$ is the equilibrium concentration at the interface (water / NAPL or water / solid) and C_w is the concentration in the bulk water at time *t*, which eventually will reach $C_{w,eq}$. The specific surface area A_o can be calculated from the porosity (n) and the degree of saturation of the pore space with NAPL (S°) yielding the volume of NAPL per unit volume of porous medium (dividing the NAPL volume by the volume of a single spherical blob of NAPL gives the number of NAPL spheres and multiplication by the surface area of one sphere finally yields the specific surface area A_0 per unit volume of porous media (in $m^2 m^{-3}$):

$$
F' = \frac{D_{aq}}{\delta} \frac{nS^o 4\pi r^2}{\frac{4}{3}\pi r^3} (C_{w,eq} - C_w)
$$

$$
F' = \frac{D_{aq}}{\delta} \frac{nS^o 3}{r_b} (C_{w,eq} - C_w) = k A_o (C_{w,eq} - C_w)
$$
 (2.10)

k is the mass transfer coefficient $(= D_{aq}/\delta [m s^{-1}])$ and *k* times A_o denotes a rate coefficient $[s^{-1}]$. In a water volume traveling through a NAPL contaminated zone, C_w increases with time (or distance at constant velocity) and finally approaches the equilibrium concentration $C_{w,eq}$.

Box 2.1: Boundary layers

The boundary layer approach to describe mass transfer between phases is based on the analogy between mass- and momentum transport in the fluid boundary layer at the interface. This approach originally has been successfully used in heat transfer (see Bennett and Myers, 2008: Momentum, Heat, and Mass Transfer, McGraw-Hill, 3rd ed.).

The boundary layer theory can be applied to the flow of fluids with high viscosity (e.g., water) at low velocity (laminar flow) through flow channels characterized by small characteristic dimensions (i.e., groundwater flow, fixed bed adsorption processes, packed-bed contactors). The boundary layer analogy presumes that the velocity and concentration profiles in the boundary layer are similar in shape, and the mass transport rate is proportional to the momentum transfer rate.

Friction/shear between the moving fluid and the surface causes the velocity of the streamlines (v_0) to decrease as we approach the surface. The resulting velocity and concentration profile are similar (right).

For estimation of the boundary layer thickness (δ) the Sherwood number is frequently used which relates the mass transfer coefficient (k) to the aqueous diffusion coefficient (D_{aa}) and size (d) :

$$
Sh = \frac{k \, d}{D_{aq}} = \frac{d}{\delta}
$$

Sh may be estimated by empirical relationships based on the Reynolds and Schmidt numbers:

$$
Sh \sim Re^{1/2} Sc^{1/3}
$$

$$
Re = \frac{v_a d}{\mu}, \quad Sc = \frac{\mu}{D_{aa}} = \frac{\delta_{hydrodynamic}}{\delta_{concentration}}
$$

Re (Reynolds number) represents the ratio of inertial to viscous forces (i.e., dynamic pressure vs. shearing stress), and it is a quantitative indicator of the amount of turbulence in the boundary layer. The boundary layer theory approximation is valid for $0 < Re < 3x10^5$ (v_a : mean or average velocity, *d*: characteristic length, e.g. grain diameter; μ : kinematic viscosity = dynamic viscosity/density, typically $1x10^{-6}$ m² s⁻¹).

Sc (Schmidt number) represents the ratio of kinematic viscosity to molecular diffusivity (D_{qq}) : aqueous diffusion coefficient). This ratio accounts for differences between the diffusivities of solutes, as well as adjusting for the differences between the thickness of the velocity and the concentration profiles (it physically relates the relative thickness of the hydrodynamic layer to the mass-transfer boundary layer).

The dissolution rate changes with time (decreasing concentration gradient, see Fig. 2.4) and introducing the water-filled pore volume yields the temporal change in concentration:

$$
\frac{d C_w}{d t} = \frac{F'}{n_e} = k \frac{A_o}{n_e} (C_{w,eq} - C_w) = k \frac{nS^o}{n_e r_b} (C_{w,eq} - C_w)
$$
\n(2.11)

 n_e is the water-filled porosity, which is the total porosity *n* minus the NAPL filled porosity $(n - n S^{\circ})$. A_o/n_e is a specific surface area per volume of pore water (sometimes denoted as A°). *n* S°/n_e is the NAPL/water volumetric ratio. If S° is low (as often the case), the difference between n_e and *n* is negligible, and A_o/n_e simply is 3 S°/r_b .

Fig. 2.3: Film diffusion and gradients during the mass transfer of a solute form a droplet of NAPL (or a mineral) across a characteristic distance δ with time into a finite bath; red, green line: initial, intermediate concentration gradient; dashed line: equilibrium conditions establish after a particular time or transport distance (and no net mass transfer occurs anymore).

In a flow-through system (groundwater or column percolation), the change of concentration with distance is relevant, and we may replace *t* by x/v_a (if the water travels with constant velocity):

$$
v\frac{d\ C_w}{d\ x} = \frac{d\ C_w}{d\ t} = k\frac{A_o}{n_e}(C_{w,eq} - C_w) = > \frac{d\ C_w}{d\ x} = \frac{k\ A_o}{v_a n_e}(C_{w,eq} - C_w)
$$
 (2.12)

which upon integration yields a simple analytical solution for the change of concentration with distance traveled:

$$
\int_{0}^{C_{w}} \frac{d C_{w}}{C_{w,eq} - C_{w}} = \int_{0}^{x} \frac{k A_{o}}{n_{e} v_{a}} dx
$$
\n
$$
-\ln(C_{w,eq} - C_{w}) + \ln(C_{w,eq}) = -\ln\left(1 - \frac{C_{w}}{C_{w,eq}}\right) = \frac{k A_{o}}{n_{e} v_{a}} x
$$
\n
$$
\frac{C_{w}}{C_{w,eq}} = 1 - \exp\left(-\frac{k A_{o}}{n_{e}} \frac{x}{v_{a}}\right)
$$
\n(2.13)

The length of mass transfer zone is defined for an argument of the exponential function of -1, and thus yields for $C_w/C_{w,eq} = 0.632 (= 1 - \exp(-1))$:

$$
X_s = \frac{v_a n_e}{k A_o} = \frac{v_a n_e r_b}{k n S^o 3}
$$
 (2.14)

This may also be easily derived from initial fluxes $(C_w = 0)$, and calculating the transport distance needed to achieve $C_{w,eq}$ in water based on eq. 2.12. If n_e is close to *n* (which often is the case as S° is often lower than 5%) eq. 2.14 reduces to:

$$
X_s \approx \frac{v_a r_b}{3 k S^o} = \frac{2 v_a r_b^2}{3 S h D_{aq} S^o}
$$
 (2.15)

If a fixed Sherwood number is used to estimate k ($k = Sh$ D_{aq} /(2 r_b)), then X_s depends on the square of the grain or blob size. Since Sh depends as well on the grain size, different relationships between the X_s and r_b may be obtained (e.g., $X_s \propto r_b^{3/2}$, see Fig. 2.2).

Note that the dissolution of minerals in a packed bed (for easily soluble calcite or gypsum crystals) may be described by the same approach using the volume of the solids (per unit volume) $(1 - n)$ and assuming the total pore volume occupied by water. X_s , in this case, is v_a *n r* /(3 *k* (1 - *n*)). X_s may be considered as the "saturation length" as used, e.g., for the description of the dissolution kinetics of carbonate sand by Schulz (1988). X_s is independent of the solubility of the compound ($C_{w,eq}$ or C_o) but varies significantly according to *k* and A_o . The definitions of *k* and A_o depend on the governing process (NAPL dissolution or desorption). For barely soluble minerals, intracrystalline processes are limiting and not the aqueous boundary layer.

At $x = X_s$, the contaminant or solute concentration in the groundwater reaches 63.2% of the equilibrium concentration $(C/C_0 = 0.632)$ as shown in Fig. 2.3. At distances less than X_{Si} , contaminant release occurs at maximum rates (F_{max}) . At larger distances $(x \gg X_s)$, the maximum contaminant concentration $(C_o = C_{w,eq})$ is reached in the groundwater.

Fig. 2.4 shows X_S calculated as a function of the groundwater flow velocity based on various empirical Sherwood correlations from the literature. Although significant differences between the published relationships are obvious, X_S is typically less than one meter in all cases, and for flow velocities below 1 m/d, even less than 10 cm. In laboratory experiments on dissolution kinetics, X_s was observed in the range of 1 cm - 2 cm (Imhoff et al., 1996), < 14 cm (Powers et al., 1994), and < 10 cm (Geller and Hunt, 1993). In heterogeneous sediments, where the distribution of the residual NAPL is highly variable, the length of the mass transfer zone may be larger than in 1D column experiments due to "dissolution fingering" (Mayer and Miller, 1996; Imhoff and Miller, 1996; Imhoff et al., 1996). In some cases, NAPLs may wet the surfaces of the grains and A_0 then is given by the surface area of the oil-wet grains. Since this is much larger than A_o for blobs in residual phase, fast dissolution kinetics are expected, resulting in short X_s .

Fig. 2.4: Increase of C/C_o in a contaminated zone. Initially, the contaminant release occurs at maximum flux. After a certain distance, the concentration gradient levels off and equilibrium conditions (maximum concentration) are reached. X_S denotes the length of the mass transfer zone.

Fig. 2.5: Comparison of different empirical correlations for the calculation of X_S based on the Sherwood number: Miller et al., 1990; Imhoff et al. (1993); Powers (1994); Fitzer et al. (1995); all calculations based on: $d = 2$ mm; $r_b = d/4$; $\theta = 0.35 \cdot 0.05 = 0.0175 \Rightarrow A_o = 105$ m² /m³; $n_e = 0.35 \cdot \theta = 0.3325$;

2.2 DISSOLUTION OF NAPL POOLS

As NAPLs infiltrate into the subsurface, they may encounter less permeable strata, which act as capillary barriers impeding further vertical movement, or they may reach the base of the aquifer. The result is the formation of a coherent lateral body of the NAPL phase, a "pool," e.g., at the bottom of the aquifer. The contaminant release rates out of NAPL pools depend on the contact time between the aqueous phase and the contaminated region. Due to the less favorable surface/volume ratio, pools dissolve slower than dispersed residual phase (blobs).

The dissolution behavior of pools can be calculated assuming that dissolution occurs due to diffusion and transverse vertical dispersion of the NAPL constituent(s) according to Fick's 2nd law.

The concentration profile which develops as the groundwater passes the pool can be calculated using the analytical solution for the semi-infinite case and a constant source (Hunt et al., 1988; Johnson and Pankow, 1992):

$$
\frac{C}{C_0} = \text{erfc}\left(\frac{z}{2\sqrt{D\frac{x}{v_a}}}\right)
$$
\n(2.16)

 x/v_a represents the residence time of the groundwater in contact with the pool. v_a is the groundwater flow velocity parallel to the pool. For a pool of length L_p , the contact time t_c equals L_p/v_a . *z* denotes the vertical distance above the pool. *D* accounts for the diffusion and transverse vertical dispersion coefficient $[L^2 s^{-1}]$:

$$
D = D_p + \alpha_t v_a \tag{2.17}
$$

 α_t and D_p denote the transverse dispersivity [L] and the pore diffusion coefficient [L² t⁻¹] in the aquifer (in a first approximation D_{aq} times porosity), respectively.

Fig. 2.6: Isolines of relative concentrations (left) and the concentration profile developed downgradient of a pool (groundwater flow velocity 1 m/d; $D_p = 1.75 \times 10^{-10} \text{ m}^2/\text{s}$; $\alpha_t = 0.5 \text{ mm}$) which close to the pool surface decreases linearly in vertical direction (right).

Fig. 2.7: Concentration profiles measured for different compounds (BTEX: left, PAHs: right) 5 cm before the end of a 2.5 m long coal tar pool compared to profiles calculated with different values for α_t and eq. 2.16; flow velocity: 1.7 m $d⁻¹$ (Eberhard and Grathwohl, 2002).

Box 2.3: Mixing cell model for dispersion dominated pool dissolution (Loyek, 1998)

The porous medium is assumed as a (3D) tetrahedral packing of spheres; flow is from left to right, and at the lower boundary (interface to the pool surface), the equilibrium concentration is applied (pool as an infinite source). Mixing is assumed as complete in the pore throats. The distance of the mixing cells (pore throats) in the flow direction (x) and perpendicular (z) to that is the one-grain radius $(r = 1)$. "Steady-state" conditions are assumed $(dC/dt = 0)$. The number of mixing cells (MZ) in the flow direction, in this case, depends on the length of the pool l_p , and the grain radius (r_a) :

Scheme of mass transfer from the pool surface assuming complete mixing in the pore throats. Numbers indicate the distribution of particles in the porous medium (at the right-hand side calculated in an Excel spreadsheet). At the interface, the equilibr<u>ium concentration ($c_{\rm i,w}^{\rm sat}$)</u> of 400 (particles) is assumed.

Displayed above are concentration profiles at the end of a pool (90 cm long) for different grain radii (r_a) in a porous medium. Symbols denote the concentrations calculated with the mixing-model (Excel); lines were calculated with the analytical solution of Ficks $2nd$ law (erfc). Note that the concentration profile is independent of the flow velocity, and α_t is a function of the grain radius (α_t = 0.5 r_a).

Pool dissolution rates. The mass, which has diffused into the groundwater while in contact with the pool is:

$$
M = 2 C_o n \sqrt{\frac{D t_c}{\pi}} L_p B_p \tag{2.18}
$$

 C_o denotes the equilibrium concentration of the solute at the interface between pool and groundwater, i.e., is the solubility of a pure compound or the equilibrium aqueous concentration according to Raoult's law for an organic mixture. B_p is the width of a representative cross-section of the pool (e.g., a central strip, 1 m wide through the center of the pool). t_c is the contact time (= L_p/v_a). The overall dissolution rate of a pool of length L_p is:

$$
F_p = 2 C_o n \sqrt{\frac{D}{\pi t_c}} L_p B_p = 2 C_o n \sqrt{\frac{D}{\pi \frac{L_p}{v_a}}} L_p B_p = 2 C_o n \sqrt{\frac{D L_p v_a}{\pi}} B_p
$$
 (2.19)

Much simpler is the calculation of the dissolution rate based on other "equivalent" models from the boundary theory, as described below.

Equivalent models. Other models that assume film diffusion (Fick's 1st law) or a groundwater layer of a given height saturated with the solute (see Fig.) may be considered as alternatives to the semi-infinite diffusion/dispersion model for the prediction of the dissolution kinetics of pools. Eq. 2.18 can be written equivalently to Fick's 1st law, allowing the calculation of the effective film thickness:

$$
F_p = D n \frac{C_o}{\delta} L_p B_p \Longrightarrow \delta = \sqrt{D t_c \frac{\pi}{4}}
$$
 (2.20)

Similarly, the thickness of a layer of water that is saturated with the solute can be calculated based on the pool dissolution rate (Eq. 2.18) and the groundwater flow rate in that layer ($Q = B_p Z_s$ v_a *n*). In essence, the "saturation height" Z_s is:

$$
C = \frac{F_p}{Q} = \frac{2 C_o n \sqrt{\frac{D L_p v_a}{\pi}} B_p}{v_a n B_p Z_s}
$$

\n
$$
\Rightarrow Z_s = \frac{2 C_o \sqrt{\frac{D L_p v_a}{\pi}}}{C_o v_a} = \sqrt{D t_c \frac{4}{\pi}}
$$
 (2.21)

In both cases δ and Z_s are close to the square root of the mean square displacement $(z^2 = D t_c)$:

$$
Z_s = \sqrt{Dt_c}\sqrt{4/\pi}; \ \delta = \sqrt{Dt_c}\sqrt{\pi/4}
$$
 (2.22)

It should be noted that z^2 (and thus also Z_s) depends on the transverse vertical dispersivity, which dominates the dissolution behavior at high flow velocities and/or high values of α_t . If $(\alpha_t \cdot v_a)$ is much higher than D_p , z^2 becomes independent of the flow velocity $(z^2 = D t_c = (D_p + \alpha_t v_a) L_p/v_a$; for α_t v_a >> D_p : $z^2 = \alpha_t L_p$ and Z_s depends only on α_t and the length of the pool.

$$
Z_s \approx \delta \approx \sqrt{\alpha_t \, L_p} \tag{2.23}
$$

If $(\alpha_t v_a)$ is large compared to the pore diffusion coefficient, the dissolution rate becomes directly proportional to the flow velocity (Fig. 2.9).

Fig. 2.8: Equivalent models for the calculation of the dissolution rates from NAPL pools (left: semi-infinite diffusion, \bar{z} = mean square displacement $(\alpha_t x)^{0.5}$; middle: Fick's 1st law, δ : film thickness (= $(\pi/4)$ α_t *x*)^{0.5}); right: saturation concentration *C*_o within a "boundary" layer of height *Z*_S = $(4/\pi \alpha_t x)^{0.5}$.

Fig. 2.9: Dependency of the pool dissolution rate F_p (2 m long and 1 m wide) on the flow velocity with different values of transverse vertical dispersivity (α_t) ; $D_p = 1.75 \times 10^{-10}$ m² s⁻¹; $C_o = 1$ mg L⁻¹ (at higher solubilities the dissolution rates increase correspondingly).

Box 2.4: Transverse dispersion coefficients - dispersivities

The local transverse dispersivity α_t is much smaller than the longitudinal dispersivity (which to a large extent comes from differential advection) and more than one order of magnitude smaller than the mean grain size. The diagram shows vertical transverse dispersivities α_t at various flow velocities. A decrease with increasing flow velocity occurs, indicating incomplete mixing at pore throats (all data were obtained in fully saturated media, except the data from Susset (1998) and Klenk (2000) which result from transport across the capillary fringe (Klenk and Grathwohl, 2002). The apparent decline of transverse dispersivities leads to a nonlinear empirical correlation to estimate the transverse dispersion coefficient based on the Peclet number (Pe):

$$
Pe = \frac{v_a d}{D_{aq}}
$$

 D_t D_{aq} $= \frac{3}{2} (\frac{4}{\pi})$ $\frac{1}{\pi}$ Pe) 0.5 (mass transfer across the capillary fringe: Klenk and Grathwohl, 2001) D_t D_{aq} $=\frac{D_p}{P}$ D_{aq} $+$ 0.28(Pe)^{0.72} (tracer experiments fully saturated: Olsson and Grathwohl, 2006) D_t D_{aq} $=\frac{D_p}{D}$ D_{aq} $+\sqrt{\frac{Pe^2}{Pe+2+4+5.5^2}} \Rightarrow D_t = D_p + v_a \frac{d}{\sqrt{Pe+123}}$ (Chiogna et al., 2010)

Comparison of empirical equations for prediction of D_t based on the Peclet number (Pe):

Box 2.5: Film thicknesses

Note the similarity of the relationships obtained for pool dissolution with empirical correlations to estimate Sherwood numbers, e.g., from Liu et al. (2014) who report $Sh = 0.1$ (Pe)^{1/2} for pore-scale dissolution of residual NAPL:

$$
\delta = \frac{d}{Sh} = \frac{d}{0.1\sqrt{Pe}} = \frac{d}{0.1\sqrt{\frac{v}{D_{aq}}}} = 10\sqrt{\frac{D_{aq}d}{v_a}}
$$

This is similar to the expression we get for the boundary layer thickness in pool dissolution, e.g., eq. 2.21, assuming that at the pore-scale D_t is replaced by D_{aq} and the length of the pool represents the grain size:

$$
\delta = \sqrt{\frac{\pi}{4} \frac{D_t L_{pool}}{v_a}} \cong \sqrt{\frac{D_{aq} d}{v_a}}
$$

The discrepancy of these two approaches (i.e., factor 10) reflects the uncertainty in using empirical relationships for different scenarios (pore-scale *vs.* large scale). If we take, e.g., the classical Fitzer relationship, then the film thickness becomes:

$$
\delta = \frac{d}{Sh} = \frac{d}{1.9Re^{1/2}Sc^{1/3}} = \frac{d}{1.9\left(\frac{dv_a}{\mu}\right)^{1/2}\left(\frac{\mu}{D_{aq}}\right)^{1/3}} = \frac{\sqrt{\frac{d\mu}{v_a}}}{1.9\left(\frac{\mu}{D_{aq}}\right)^{1/3}} \approx 0.5\sqrt{\frac{D_{aq}d}{v_a}}
$$

The approximation is obtained by assuming an empirical exponent of the Schmidt number of $\frac{1}{2}$ (note, the product of $Re \times Sc$ is the Peclet number ($Pe = v_a d/D_{aq}$). This is similar to pool dissolution but yields different film thicknesses compared to Liu et al. (2014), as we saw already in Fig. 2.5 comparing different empirical relationships.

3. ASSESSMENT OF NAPL CONTAMINATED SITES

3.1 DEPTH-AVERAGED CONCENTRATIONS

The most important conclusion for risk assessment from the above discussion is that, in areas which are contaminated by residual NAPL (blobs trapped in the porous media, e.g., a smear zone of oil in the capillary fringe), equilibrium is reached after a short flow distance and the contaminant concentration in the "contact" groundwater is given by the water solubility of the NAPL (for organic mixtures according to Raoult's law, see next session). The depth-averaged contaminant concentration downgradient from a contaminant source is then simply given by the dilution rate. In essence, the thickness of the contaminated zone Z_0 compared to the thickness of the aquifer *h*:

$$
C_{avg} = C_o \frac{z_o}{h} \tag{3.1}
$$

For the pool dissolution scenario, z_0 is given by the height of the chemical boundary layer Z_s .

3.2 RETARDATION OF CLEAN WATER FRONTS IN SMEAR ZONES

Remediation of subsurface contamination by in situ techniques, such as pump-and-treat methods in the saturated zone or vapor phase extraction for removal of volatile compounds from the unsaturated zone, is most efficient if equilibrium conditions are achieved - high flow rates of water or air result in fast removal of the contaminants. If equilibrium conditions apply for the dissolution of residual NAPLs ("blobs"), then the time scale for remediation can be estimated from the velocity of the dissolution front (v_{Dis}) , which is retarded compared to the groundwater flow velocity (v_a) . The retardation factor is given by the ratio of the overall mass of contaminant present in the porous medium to the mobile (i.e., dissolved) contaminant:

$$
R = \frac{v_a}{v_{Dis}} = \frac{n_e C_o + \rho_o n S^o}{C_o n_e} \tag{3.2}
$$

This retardation factor equals the number of pore volumes, which have to be replaced until the complete dissolution of the NAPL is achieved. ρ_0 n S° represents the NAPL mass per unit volume soil, and $n_e C_0$ is the mobile fraction of the contaminant (i.e., dissolved in pore water). ρ_0 denotes the density of the organic phase (NAPL, the same unit as C_0 e.g., kg m⁻³), *n* is the porosity, S° is the degree of saturation with NAPL) and n_e is effective porosity $(= n - n S^{\circ})$. Since the solute mass dissolved in the pore water is, in most cases, insignificant compared to the mass of the NAPL, Eq. 3.2 reduces to:

$$
R = \frac{\rho_o n S^o}{C_o n_e} = \frac{\rho_o S^o}{C_o} \tag{3.3}
$$

If n_e is close to the overall porosity (*n*) Eq. 3.3 becomes simply: $R = \rho_o S^{\circ}/C_o$ (S[°] is the degree of saturation). As long as equilibrium conditions apply, high flow velocities (e.g., close to a well during application of pump-and-treat techniques) will result in fast dissolution/remediation of residual NAPL (in contrary to non-equilibrium, when diffusion is limiting, then release rates are independent of the groundwater flow velocity).

3.3 TIME FOR DISSOLUTION OF POOLS

The dissolution of NAPL from pools is much slower than the equilibrium dissolution of residual blobs of NAPL (this is due to the much lower interfacial area between NAPL and the mobile water). The time for pool dissolution can easily be calculated from the dissolution rate F_p (based on Z_s) and the mass of NAPL in the pool (assuming that the pool shrinks in height but not in length during dissolution and considering only a rectangular strip in the center of the pool, e.g., 1 m width):

$$
t_{Dis} = \frac{M_p}{F_p} = \frac{L_p H_p \rho_o n S^o}{C_o v_a n Z_s}
$$
\n
$$
(3.4)
$$

 H_P denotes the pool height (the NAPL saturation S° in the pool is close to 1). Fig. 3.1 shows some example calculations with a value for α_t of 0.1 mm. Even at flow velocities of 1 m/d, the time scale of TCE dissolution from pools is of the order of several decades to a century. Compounds with lower solubilities (C_0) would need even more time to be dissolved from a pool. At slow flow velocities, the dissolution time depends on the square root of v_a , whereas at higher flow velocities, the time for pool dissolution decreases linearly with increasing v_a .

Fig. 3.1: Time scales for dissolution of TCE pools (solid line: pool length = 2 m, 100 kg TCE; dotted line: pool length = 16 m, 800 kg TCE; other lines: pool lengths = 4 m and 8 m, 200 kg and 400 kg TCE, respectively; $D_p = 2E - 10 \text{ m}^2 \text{ s}^{-1}$; $\alpha_t = 0.1 \text{ mm}$).

Box 3.1: "Stacked" NAPL-pools in heterogeneous porous media

NAPL in the subsurface are usually not uniformly distributed, which means that contaminant concentrations in water show a high special variability. Water in areas containing residual NAPL is near the saturation concentration, while water outside shows low concentrations. Downstream water mixes through transverse dispersion. Various attempts have been made to simulate dissolution from multiple pools in the subsurface. An example is depicted below.

More on dissolution and dispersion in heterogeneous porous media in:

- Ronald W. Falta, WATER RESOURCES RESEARCH, VOL. 39, NO. 12, 1360, 2003: Modeling sub-gridblock-scale dense nonaqueous phase liquid (DNAPL) pool dissolution using a dual-domain approach
- Werth, C. J., O. A. Cirpka, and P. Grathwohl (2006): Enhanced mixing and reaction through flow focusing in heterogeneous porous media, Water Resour. Res., 42 (12), W12414, doi:10.1029/2005WR004511

3.4 TIME FOR DISSOLUTION OF SINGLE SPHERES (OIL DROPLET, GAS BUBBLE, CRYSTAL, ETC.)

The change in mass (*M*) of a sphere during dissolution (e.g., sugar crystal in tea) depends on the mass transfer coefficient (k) , the saturation concentration at the interface (C) ; often corresponds to the solubility of the material), and the interfacial area (the surface area of a sphere with radius *r* is: $A = 4 \pi r^2$:

$$
\frac{dM}{dt} = -k \, C \, A
$$
\n
$$
\frac{dM}{dt} = -k \, C \, 4 \, \pi r^2 \text{ and with } r = \left(\frac{3M}{4\pi\rho}\right)^{1/3}
$$
\n
$$
\frac{dM}{dt} = -k \, C \, 4 \, \pi \left(\frac{3 \, M}{4 \, \pi \, \rho}\right)^{2/3}
$$
\n(3.5)

p denotes the density of the sphere (its "concentration per volume"). The product of *k C* represents the surface area normalized dissolution rate (e.g. in g m⁻² s⁻¹, if multiplied with the molecular weight in mol m^{-2} s⁻¹), which typically is reported in literature for the dissolution of minerals. Integration from the initial radius r_0 to r and for the time from 0 to t yields:

$$
\int_{M_0}^{M} \frac{d M}{M^{2/3}} = \int_{o}^{t} -kC 4\pi \left(\frac{3}{4\pi \rho}\right)^{2/3} dt
$$

$$
3M_o^{1/3} - 3M^{1/3} = 0 + k C 4 \pi \left(\frac{3}{4 \pi \rho}\right)^{2/3} t
$$

$$
M^{1/3} = M_o^{1/3} - k C \frac{4}{3} \pi \left(\frac{3}{4 \pi \rho}\right)^{2/3} t
$$
 (3.6)

with *M* = 4/3 π r^3 ρ we get the solution for the radius as a function of time:

$$
r = r_o - \frac{k C}{\rho} t \tag{3.7}
$$

From this, the time needed to dissolve the sphere $(r = 0)$ can be calculated by:

$$
t_{diss} = \frac{r_o \rho}{k C} \tag{3.8}
$$

For comparison, the time needed to dissolve the sphere based on the maximum initial flux (without accounting for the shrinking surface area of the sphere) is three times shorter:

$$
\frac{4}{3}\pi r^3 = k\mathcal{C} 4\pi r^2 t \implies t_{diss} = \frac{r_o \rho}{3\ k\ \mathcal{C}} \left(= \frac{M_p}{F_p} \right) \tag{3.9}
$$

Note the similarity to eq. 3.4.

Accounting for *k* **– size (e.g.,** *d* **= 2** *r***) relationships. The mass transfer coefficient (***k***) depends** on the radius (r) in empirical relationships such as the Sherwood number (Sh) :

$$
k = \frac{D_{aq}}{\delta} \text{ and } Sh = \frac{d}{\delta} = \frac{2 r}{\delta}
$$

$$
\implies k = \frac{D_{aq} Sh}{2 r}
$$
 (3.10)

Solving the mass balance equation directly for the radius we get:

$$
d(M) = (-k C 4 \pi r^2)dt
$$

\n
$$
d\left(\frac{4}{3} \pi r^3 \rho\right) = \left(-\frac{D_{aq} Sh}{2 r} C 4 \pi r^2\right) dt
$$

\n
$$
\frac{1}{r} dr^3 = \left(-D_{aq} \frac{3}{2} Sh \frac{C}{\rho}\right) dt
$$
\n(3.11)

Integration (with $r^3 = a$, $r = a^{1/3}$, $r^2 = a^{2/3}$ and re-substitution) yields:

$$
\int_{r_0}^r \frac{1}{r} dr^3 = \int_0^t \left(-D_{aq} \frac{3}{2} Sh \frac{C}{\rho} \right) dt
$$

$$
r_0^2 - r^2 = D_{aq} Sh \frac{C}{\rho} t
$$

$$
r = \sqrt{r_0^2 - D_{aq} Sh \frac{C}{\rho} t}
$$
 (3.12)

The time to dissolve the sphere in this case is:

$$
t_{diss} = \frac{r_o^2 \rho}{D_{aq} Sh C}
$$
 (3.13)

Thus, at a given Sherwood number (often between 2 and 6), the time to dissolve the sphere increases with the radius squared. Since *Sh* often is found to be a function of the square root of Re or *Pe*, the dissolution time scales with the radius to the power of 3/2.

4. DISSOLUTION OF MIXTURES

4.1 RAOULT'S LAW AND PARTITION COEFFICIENTS

As shown above, the driving force for mass transfer is the concentration at the interface, which depends on the partition coefficient of compound *i* between the organic mixture and water phase. In complex mixtures of organic compounds (e.g., coal tar, gasoline), the aqueous phase saturation concentration of the individual components $C_{i,w}$ depends on the composition of the mixture (Banerjee, 1984; Mackay et al., 1991; Lane and Loehr, 1992; Lee et al., 1992). $C_{i,w}$ is always less than the solubility of the pure substances in water and can be determined in liquid/liquid mixtures following Raoult's Law (Pyka, 1994; Loyek and Grathwohl, 1995):

$$
C_{i,w} = \chi_{i,o} \gamma_{i,o} S_i \tag{4.1}
$$

where $C_{i,w}$, $\chi_{i,o}$ and $\chi_{i,o}$ denote the concentration of component *i* in water [g l⁻¹ or mol l⁻¹]; the molar fraction of component *i* in the organic mixture $\lceil n_i / \sum n_i \rceil$; and the activity coefficient of *i* in the organic mixture (most =1); S_i is the solubility in water of the individual component *i* (pure substance) in g $l¹$ or mol $l¹$ (for solids of the subcooled liquid, for gases of superheated liquid; or the enhanced solubility in the presence of cosolvents (e.g., alcohols).

The activity coefficient $(\gamma_{i,o})$ describes the deviation from the ideal behavior. For an ideal mixture: $\gamma_{i,o}$ = 1, meaning Eq. 4.1 simplifies to:

$$
C_{i,w} = \chi_{i,o} S_i \tag{4.2}
$$

The molar fraction $\chi_{i,o}$ can have a maximum value of 1 for a pure compound, and can be easily calculated from the mean molecular weight of the mixture $(M_0 \, [\text{g mol-1}])$, the molecular weight of *i* [g mol⁻¹] and the fraction of *i* within the mixture $(f_{i,o}$ or weight-%/100):

$$
\chi_{i,o} = f_{i,o} \frac{M_o}{M_i}; \text{for } M_o \cong M_i \text{ then } \chi_{i,o} \cong f_{i,o}
$$
 (4.3)

In many cases, M_o/M_i in organic liquids (e.g., fuels) is between 0.5 and 1.5. For the solubility from organic mixtures, the following equation applies:

$$
C_{i,w} = f_i \frac{M_o}{M_i} S_i
$$

and in the first rough approximation $(M_o/M_i \cong 1)$: (4.4)

$$
C_{i,w} = \frac{weight\%_i}{100} S_i
$$

 $C_{i,w}$ is always less than the solubility of the pure substance (also when the solubility of the subcooled liquid is higher than the solubility of the solid substance: precipitation of crystals would occur if the solubility of the pure substance is exceeded).

4.1 DOUBLE FILM DIFFUSION

In NAPL mixtures (e.g., fuels) mass transfer of a constituent may be limited by diffusion in the NAPL and water (or gas phase during volatilization). The double-film diffusion model approximates the NAPL-water interface by introducing two films adjacent to the interface, where one film represents a stagnant water layer and a second layer in the multi-component organic phase (see Fig. 4.1). The concentration of contaminants across this multi-component organic film decreases from the initial concentrations (C_{i_0}) in the organic phase to the concentration at the NAPL-water interface $(C_{i_0/w})$. At the boundary layer between the two films, organic contaminants leave the organic phase and enter the aqueous phase. The concentrations at the NAPL-water interface are at equilibrium and given by the organic-water partition coefficient $(K_{o/w})$ of the respective compounds.

Box 4.1: Partition coefficients between organic mixtures and water

Since $f_{i,o}$ represents the concentration of *i* in the organic phase mixture (in kg kg⁻¹), the partitioning coefficient between an organic phase and water (K_0/w) can be derived from Raoult's Law using:

$$
K_{o/w} = \frac{C_{i,o}}{C_{i,w}} = \frac{f_{i,o}}{f_{i,o} \frac{M_o}{M_i} S_i} = \frac{1}{\frac{M_o}{M_i} S_i}
$$

The partition coefficient is independent of concentration. For $M_0 = M_i$ or a pure solvent (e.g., Toluene in Toluene: $C_{i,o}$ = 1 kg kg⁻¹; $C_{i,w}$ = S_i). The partitioning coefficient between a pure organic phase ("solvent") and water then would be simply: $K_{o/w} \approx \frac{1}{S_i}$ (= "ideal" case)

Box 4.2: Retardation factor of the "clean waterfront" moving through a smear zone of residual phase during the dissolution of a compound form an organic NAPL mixture

How fast the dissolution of a NAPL constituent happens depends on the retardation factor. The retardation factor is defined as:

$$
R_d = \frac{total\ concentration}{mobile\ concentration} = \frac{C_{i,w} n_e + n S^o \rho_o C_{i,o}}{C_{i,w} n_e} = 1 + \frac{C_{i,o} n S^o \rho_o}{C_{i,w} n_e}
$$

n S° represents the oil-filled pore volume [e.g., m³ m⁻³], (*n* S° ρ_o) the mass of oil [kg m⁻³] and (*n* S° ρ_o $C_{i,o}$) finally the mass of the compound in the NAPL per volume of porous media [kg of *i* per m³]. n_e is the effective water-filled pore volume = $n - n S^o$ (the pore volume minus the NAPL volume). $C_{i,o}/C_{i,w}$ represents the partition coefficient $K_{o/w}$ (which comes from Raoult's law, as shown in Box 1.1). *n* S° ρ_o $/n_e$ is oil to water ratio in the porous media. For a pure NAPL (e.g., pure toluene) $C_{i,w}$ is represented by the water solubility of the compound S_i . Thus, the retardation factor is:

$$
R_d = 1 + K_{o/w} \frac{n S^o \rho_o}{n_e} \approx 1 + \frac{n S^o \rho_o}{\frac{M_o}{M_i} S_i n_e}
$$

As $K_{o/w}$, R_d is independent of the composition of the organic mixture. The term to the right is obtained by implementing Raoult's law with an activity coefficient of 1. Thus, the movement of a clean waterfront through a NAPL smear zone is independent of the solute content in the NAPL, but depends on the amount of NAPL in the porous medium. The equation further simplifies if the molecular weights of the NAPL and compound *i* are assumed to be similar $(M_0 \cong M_i)$ and $n_e \cong n$:

$$
R_d \cong 1 + \frac{nS^o \rho_o}{S_i n_e} \cong 1 + \frac{S^o \rho_o}{S_i}
$$

Fig. 4.1: Schematic view of the stagnant two-film model. Mass transfer across both films is assumed to be controlled by Fick's 1st law. At the interface between both films, "local equilibrium" is assumed, and concentrations on either side are controlled by the partition coefficients of the component between the organic phase and water. δ_o and δ_w refer to the respective film thicknesses; $C_{i,o}$ and $C_{i,w}$ are the concentrations in the organic and aqueous phase, respectively.

To calculate the contaminant flux across the two films, the following assumptions have to be made:

- \circ the concentration in the mobile water is constant (e.g. close to zero if water continuously displaced during remediation): $C_{i,w} = 0$
- \circ steady-state conditions apply for the flux of contaminant across both layers (no storage)
- \circ the concentration at the NAPL-water interface $(C_{i,o/w})$ is at equilibrium with the concentration at the water-organic interface $(C_{i,w/o})$ (e.g., saturation concentration in water, according to Raoult's law)

Then, the fluxes in each film can be evaluated using Fick's 1st law, which for the organic film is:

$$
F_{i,o} = \frac{-D_o}{\delta_o} \Delta C_{i,o} = \frac{-D_o}{\delta_o} (C_{i,o/w} - C_{i,o})
$$
\n(4.5)

and analogously in the aqueous film:

$$
F_{i,w} = \frac{-D_{aq}}{\delta_w} \Delta C_{i,w} = \frac{-D_{aq}}{\delta_w} (C_{i,w} - C_{i,w/o})
$$
\n(4.6)

where $C_{i,o/w}$, and $C_{i,w/o}$ denote the interfacial concentrations of *i* in the organic and the aqueous phase, respectively; D_o is the diffusion coefficient in the organic mixture.

Under steady-state conditions, the fluxes in both films are equal. By introducing the partition coefficient of the compound between the organic and aqueous phase $(K_{o/w})$, the total flux can be calculated:

$$
K_{o/w} = \frac{C_{i,o/w}}{C_{i,w/o}}
$$
 (4.7)

By combining Eqs. 4.5 to 4.7 the flux is:

$$
F_{i} = F_{i,w} = F_{i,0} = \frac{C_{i,o} - C_{i,w}K_{o/w}}{\frac{\delta_{w}K_{o/w}}{D_{aq}} + \frac{\delta_{o}}{D_{o}}} = \frac{\frac{C_{i,o}}{K_{o/w}} - C_{i,w}}{\frac{\delta_{w}}{D_{aq}} + \frac{\delta_{o}}{D_{o}K_{o/w}}}
$$
(4.8)

If both films have similar thickness, then the flux will is controlled by $K_{o/w}/D_{aq}$ and $1/D_o$. Since for hydrophobic organic compounds $K_{o/w}$ is much larger than 1, the flux will be controlled by the water film. Thus, the contaminant flux can be approximated using:

$$
F_i = \frac{\frac{C_{i,o}}{K_{o/w}} - C_{i,w}}{\frac{\delta_w}{D_{aq}}}
$$
\n(4.9)

If the concentrations in the mobile phase are zero (or close to it), then Eq. 1.9 can be further simplified to:

$$
F_i = \frac{D_{aq}}{\delta_w} \frac{C_{i,o}}{K_{o/w}}
$$
(4.10)

where the ratio D_{aq} / δ_w represents the mass transfer coefficient (*k* [m s⁻¹]); $C_{i,o}$ / $K_{o/w}$ denotes the equilibrium concentration of a compound in water $(e.g.$ its water solubility) and according to Eq. 4.10, the total flux depends linearly on this aqueous concentration at the interface. A change in flow velocity or temperature will affect δ and D_{aq} equally for all compounds involved, resulting in a simultaneous change in the dissolution rate. If the interfacial concentration is increased by adding a surfactant for solubilization, contaminants with low water solubilities are affected the most, as described in the next section.

Fig. 4.2 shows experimental results of the dissolution of PAHs from a coal tar "pool." With increasing flow velocities, fluxes of all PAHs increase simultaneously, indicating that the water film thickness (δ_w) decreases as flow rates increase. Fig. 4.3 illustrates that the observed dissolution rates for a given flow velocity in this experiment are linearly proportional to the respective equilibrium aqueous phase concentrations of the individual PAHs $(= C_{i,o} / K_{o/w})$ as predicted from the film diffusion model. Since D_{aq} , $C_{i,o}$, and $K_{o/w}$ are known from such type of experiments, the film thicknesses (δ_w) can be calculated. This approach is also valid for solid organic phases such as polymers (e.g., microplastics, passive samplers).

Fig. 4.2: Influence of flow rate (velocity), temperature and surfactants on the dissolution rate of PAHs from a DNAPL (creosote) pool; changes in flow velocity and temperature affect all compound the same way whereas addition of surfactants increases dissolution rates of low solubility compounds most (because of high partitioning into micelles) (Eberhardt, 1995).

Fig. 4.3: Dissolution rates of PAHs from coal tar and creosote for a given flow velocity are linearly proportional to the respective equilibrium concentration.

5. REMEDIATION: SOLUBILIZATION BY SURFACTANTS

5.1 ENHANCED SOLUBILITY: MICELLES

The dissolution of NAPL can be accelerated if the aqueous concentration at the interface and bulk water are increased by the application of surfactants. Surfactants form aggregates in water - socalled micelles - which can "carry" hydrophobic organic compounds in their interior (see Fig. 5.1). On solid surfaces, they form bilayers (like a membrane) called admicelles. For organic mixtures, surfactants decrease the partition coefficient, which in turn increases the aqueous concentration at the interface. This process is called solubilization. In contrast, mobilization refers to the movement of the NAPL as a separate phase. Micelles form above a certain concentration of surfactants in water, which is known as the critical micelle concentration (CMC - often between 0.1 and 1 g l -1). The relative increase in aqueous solubility (denoted here as S^*) is dependent on the concentration of the surfactant monomers $(C_{mono}$ [kg L⁻¹]), the concentration of micelles $(C_{mic}$ [kg L⁻¹]), and the respective partitioning coefficients K_{mono} and K_{mic} [L kg-1]:

$$
S^* = \frac{C_{i,w}}{S_i} = \frac{C_{i,w}}{C_{i,w}^{sat}} = 1 + C_{mono} K_{mono} + C_{mic} K_{mic}
$$

\n
$$
\implies \text{for } C_{mic} \gg CMC:
$$

\n
$$
S^* = 1 + C_{mic} K_{mic}
$$
 (5.1)

Fig. 5.2 shows how "solubilization" affects low solubility compounds (= compounds with high K_{OW} values) most. Thus, surfactants "level out" the "solubility" of organic compounds, as shown in Fig. 5.3. Consequently, solubilization leads to faster dissolution, as shown in Fig. 5.4. Solubilization by surfactant micelles is equivalent to adding humic substances or colloidal matter to water, which also take up organic compounds in their hydrophobic moieties.

Fig. 5.1: Examples of common surfactants and steps of structure formation for micelles and admicelles.

Fig. 5.2: Examples for relative "solubility" enhancement for different organic compounds (S*: solubility **enhancement** factor \cong 1 + C_{mic} K_{mic})

Triton X-100 (CMC = 150 mg/l)

Fig. 5.3: Apparent increases in the solubility in water of the pure substance $s^*_{i,w}$ in mg l ⁻¹ for chosen PAHs in relation to the surfactant concentration above CMC (Triton X-100: commercial surfactant)

Release rates without surfactants

Fig. 5.4: Release rates of Nap, Phe, and BaA from coal tar in residual phase with and without surfactants (calculated after Loyek and Grathwohl, 1996): With the application of surfactants, all PAHs are more or less simultaneously dissolved.

5.2 SORPTION OF SURFACTANTS AND FACILITATED TRANSPORT BY MICELLES

When surfactants are introduced in the subsurface environment, they will as other compounds interact with the solid phase (e.g., minerals). Depending on their surface charge, surfactants adsorb on solid surfaces until the critical micelle concentration is reached (Fig. 5.5). Then a maximum loading of the surface $(q_{max}$ or $C_{s,max}$ with the surfactant is reached and a bilayer may form (admicelles). This sorption behavior is typically described by Langmuir isotherms. Since the surface area is inversely related to grain size, adsorption increases with increasing surface to volume ratio (for spheres $= 3/r$ at shown in Fig. 5.5. If sorption is known, the amount of surfactant lost during subsurface remediation can be calculated.

Fig. 5.5: Adsorption of surfactants to different aquifer materials (Terrasurf G50, grain size $2 - 4$ mm) and decreasing maximum loading (q_{max}) as a function of the surface to volume ratio (from Danzer and Grathwohl, 1998); sandstones show high adsorption of surfactants because of the internal surface area is accessible to surfactants.

Surfactants adsorbed on solid surfaces form organic coatings (admicelles) which take up organic compounds from water and this causes their retardation during transport. Above the CMC mobile micelles appear in the mobile phase (water) which in turn facilitate transport of organic compounds. Such systems may be used in two ways: Below the CMC, surfactants from "sorptive" barriers and above they facilitate removal of organic pollutants. This was also proposed for humic substances since they behave similarly to surfactants. This behavior is very simply described by the retardation factor:

$$
R_d = \frac{\text{total concentration}}{\text{mobile concentration}} \\
= \frac{C_w n + K_{mic} C_{mic} C_w n + K_{oc} f_{oc} C_w \rho_b + K_{adm} f_{adm} C_w \rho_b}{C_w n + K_{mic} C_{mic} C_w n} \\
= 1 + \frac{K_{oc} f_{oc} \rho_b + K_{adm} f_{adm} \rho_b}{n + K_{mic} C_{mic} n} = 1 + \frac{K_{oc} f_{oc} + K_{adm} f_{adm} \rho_b}{n + K_{mic} C_{mic} n} \\
\text{for } C_{mic} = 0 \tag{5.2}
$$
\n
$$
R_d = 1 + (K_{oc} f_{oc} + K_{adm} f_{adm}) \frac{\rho_b}{n} \text{ (increases with } f_{adm}) \\
\text{for } K_{mic} C_{mic} \gg 1 \text{ and } K_{adm} f_{adm} \gg K_{oc} f_{oc} \\
\Rightarrow R_d = 1 + \frac{K_{adm} f_{adm} \rho_b}{K_{mic} C_{mic} n} \approx 1 + 2 \frac{f_{adm} \rho_b}{C_{mic} n} \text{ (decreases with } C_{mic})
$$

 f_{adm} denotes the fraction of surfactant on the surface analogous to f_{oc} (fraction of organic carbon). *C_{mic}* is the micelle concentration in water (e.g., kg l⁻¹) and *K_{adm}* denotes the partition coefficient of the solute between water and admicelles. $K_{oc} f_{oc}$ is the distribution coefficient K_d ($K_{oc} f_{oc}$ C_w and K_{adm} *fadm* C_w represent the sorbed or co-sorbed concentration of the compound). Fig. 5.6 illustrates how "solubilization" affects the transport of contaminants - below the CMC the retardation factor increase and above the CMC it decreases again.

Fig. 5.6: Increase and decrease of the retardation factor for organic compounds with increasing surfactant (Terrasurf G50: ethoxylated fatty alcohol) concentration (from Danzer and Grathwohl, 1998); CMC: Critical Micelle Concentration; measured data from column experiments with phenanthrene in alluvial sand (Neckar): $K_d = 1.2$ l kg⁻¹, $K_{admicelle} = 18300$ l kg⁻¹, $K_{d,surf} = 2.34$ l kg⁻¹.

6. MORE LITERATURE

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APPENDIX

Physicochemical Properties of Gasoline Hydrocarbons

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Saturation Concentrations of Hydrocarbons from Kerosine

CONTAMINATED SITES (GRATHWOHL) NAPL DISSOLUTION (MASS TRANSFER)

Saturation Concentrations of Hydrocarbons from Kerosine (cont.)

CONTAMINATED SITES (GRATHWOHL) NAPL DISSOLUTION (MASS TRANSFER)

Saturation Concentrations of Hydrocarbons from Diesel Fuel

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 $S^{\text{w}}_{\text{i}\text{ 'sel}}$ and P_{L} are the subcooled solubility and liquid-vapor pressure in grey color, respectively.